# A Poincaré-Hopf type formula for Chern character numbers 

Huitao Feng*, Weiping Li and Weiping Zhang ${ }^{\dagger}$


#### Abstract

For two complex vector bundles admitting a homomorphism with isolated singularities between them, we establish a Poincaré-Hopf type formula for the difference of the Chern character numbers of these two vector bundles. As a consequence, we extend the original Poincaré-Hopf index formula to the case of complex vector fields.


## 1 Introduction and the statement of the main result

Let $M$ be a closed, oriented, smooth manifold of dimension $2 n$. Let $E_{+}, E_{-}$be two complex vector bundles over $M$.

Let $v \in \Gamma\left(\operatorname{Hom}\left(E_{+}, E_{-}\right)\right)$be a homomorphism between $E_{+}$and $E_{-}$. Let $Z(v)$ denote the set of the points at which $v$ is singular (that is, not invertible). We assume that the following basic assumption holds.

Basic Assumption 1.0. The point set $Z(v)$ consists of a finite number of points in $M$.
For any $p \in Z(v)$, we choose a small open ball $B(p)$ centered at $p$ such that the closure $\overline{B(p)}$ contains no points in $Z(v) \backslash p$. Then, when restricted to the boundary $\partial B(p)$, the linear map

$$
\begin{equation*}
\left.v\right|_{\partial B(p)}:\left.\left.E_{+}\right|_{\partial B(p)} \rightarrow E_{-}\right|_{\partial B(p)}, \tag{1.1}
\end{equation*}
$$

which we denote by $v_{p}$, is invertible. The map $v_{p}$ determines an element in $K^{1}\left(S^{2 n-1}\right)=\mathbf{Z}$ which we denote by $\operatorname{deg}\left(v_{p}\right) \in \mathbf{Z} .{ }^{1}$

The main result in this paper is the following theorem:
Theorem 1.1. Under the Basic Assumption 1.0, the following identity holds,

$$
\begin{equation*}
\left\langle\operatorname{ch}(E+)-\operatorname{ch}\left(E_{-}\right),[M]\right\rangle=(-1)^{n-1} \sum_{p \in Z(v)} \operatorname{deg}\left(v_{p}\right) . \tag{1.2}
\end{equation*}
$$

[^0]Our original motivation is to establish an extension of the Poincaré-Hopf index formula for vector fields with isolated zero points (cf. [1, Theorem 11.25]) to the case of complex vector fields, under the framework considered by Jacobowitz in [3].

To be more precise, let $T_{\mathbf{C}} M=T M \otimes \mathbf{C}$ denote the complexification of the tangent vector bundle $T M$. Let $K=\xi+\sqrt{-1} \eta \in \Gamma\left(T_{\mathbf{C}} M\right)$ be a smooth section of $T_{\mathbf{C}} M$, with $\xi, \eta \in \Gamma(T M)$.

Let $g^{T M}$ be a Riemannian metric on $T M$, then it induces canonically a complex symmetric bilinear form $h^{T_{\mathbf{C}} M}$ on $T_{\mathbf{C}} M$, such that

$$
\begin{equation*}
h^{T_{\mathbf{C}} M}(K, K)=|\xi|_{g^{T M}}^{2}-|\eta|_{g^{T M}}^{2}+2 \sqrt{-1}\langle\xi, \eta\rangle_{g^{T M}} . \tag{1.3}
\end{equation*}
$$

Jacobowitz proved in [3] the following vanishing result.
Proposition 1.2. (Jacobowitz [3]) If $h^{T_{\mathbf{C}} M}(K, K)$ is nowhere zero on $M$, then the Euler number of $M$ vanishes: $\chi(M)=0$.

If one takes $\eta=0$, then Proposition 1.2 reduces to the classical Hopf vanishing result: $\chi(M)=0$ if $M$ admits a nowhere zero vector field.

Jacobowitz asked in [3] whether there is a counting formula for $\chi(M)$ of Poincaré-Hopf type, extending Proposition 1.2 to the case where $h^{T_{\mathbf{C}} M}(K, K)$ vanishes somewhere on M. In Section 3, we will establish such a formula as an application of Theorem 1.1, while Theorem 1.1 itself will be proved in Section 2.

## 2 A Proof of Theorem 1.1

We will use the superconnection formalism developed in [5] to prove Theorem 1.1.
Due to the topological nature of both sides of (1.2), we first make some simplifying assumptions on the metrics and connections near the set of singularities $Z(v)$.

First of all, we assume that there is a Riemannian metric $g^{T M}$ on $T M$ such that for any $p \in Z(v)$, there is a coordinate system $\left(x_{1}, \cdots, x_{2 n}\right)$, with $0 \leq x_{i} \leq 1$ for $1 \leq i \leq 2 n$, centered around $p$ such that

$$
\begin{equation*}
B_{p}(1)=\left\{\left(x_{1}, \ldots, x_{2 n}\right) \mid \sum_{i=1}^{2 n} x_{i}^{2} \leq 1\right\} \subset M \backslash(Z(v) \backslash\{p\}) \tag{2.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\left.g^{T M}\right|_{B_{p}(1)}=d x_{1}^{2}+d x_{2}^{2}+\cdots+d x_{2 n}^{2} \tag{2.2}
\end{equation*}
$$

that is, the metric $g^{T M}$ is Euclidean on each $B_{p}(1), p \in Z(v)$.
On the other hand, on each $B_{p}(1)$, the bundles $E_{ \pm}$are trivial vector bundles. We equip these two trivial vector bundles over $B_{p}(1)$ the trivial metrics and trivial connections respectively. Moreover, we can deform $v$ near $\partial B_{p}(1)$, so that $v_{p}:\left.\left.E_{+}\right|_{\partial B_{p}(1)} \rightarrow E_{-}\right|_{\partial B_{p}(1)}$ is unitary, while still keep the new homomorphism nonsingular on $M \backslash Z(v)$.

By partition of unity, we may then construct Hermitian metrics and connections $\nabla^{E_{ \pm}}$on $E_{ \pm}$ over $M$ such that the above simplifying assumptions hold on $\cup_{p \in Z(v)} B_{p}(1)$.

We now follow the formalism in [5].
Let $E=E_{+} \oplus E_{-}$be the $\mathbf{Z}_{2}$-graded complex vector bundle over $M$. Let $\nabla^{E}=\nabla^{E_{+}} \oplus \nabla^{E_{-}}$ be the $\mathbf{Z}_{2}$-graded connection on $E$.

Let $v: E_{+} \rightarrow E_{-}$extend to an (odd) endomorphism of $E$ by acting as zero on $E_{-}$, with the notation unchanged. Let $v^{*}: E_{-} \rightarrow E_{+}$(and thus also extends to an (odd) endomorphism of $E)$ be the adjoint of $v$ with respect to the Hermitian metrics on $E_{ \pm}$respectively.

Set $V=v+v^{*}$. Then $V$ is an odd endomorphism of $E$. Moreover, $V^{2}$ is fiberwise positive over $M \backslash Z(v)$.

We fix a square root of $\sqrt{-1}$. Let $\varphi: \Omega^{*}(M) \rightarrow \Omega^{*}(M)$ be the rescaling on differential forms such that for any differential form $\alpha$ of degree $k, \varphi(\alpha)=(2 \pi \sqrt{-1})^{-\frac{k}{2}} \alpha$. The final formulas below will not depend on the choice of this square root.

For any $t \in \mathbf{R}$, let $\mathbf{A}_{t}$ be the superconnection on $E$, in the sense of Quillen [5], defined by

$$
\begin{equation*}
\mathbf{A}_{t}=\nabla^{E}+t V . \tag{2.3}
\end{equation*}
$$

Let $\operatorname{ch}\left(E, \mathbf{A}_{t}\right)$ be the associated Chern character form defined by

$$
\begin{equation*}
\operatorname{ch}\left(E, \mathbf{A}_{t}\right)=\varphi \operatorname{tr}_{s}\left[e^{-\mathbf{A}_{t}^{2}}\right] \tag{2.4}
\end{equation*}
$$

The following transgression formula has been proved in [5, (2)],

$$
\begin{equation*}
\frac{\partial \operatorname{ch}\left(E, \mathbf{A}_{t}\right)}{\partial t}=-\frac{1}{\sqrt{2 \pi \sqrt{-1}}} d \varphi \operatorname{tr}_{s}\left[V e^{-\mathbf{A}_{t}^{2}}\right] . \tag{2.5}
\end{equation*}
$$

Set for any $T>0$,

$$
\begin{equation*}
\gamma(T)=\frac{1}{\sqrt{2 \pi \sqrt{-1}}} \varphi \int_{0}^{T} \operatorname{tr}_{s}\left[V e^{-\mathbf{A}_{t}^{2}}\right] d t \tag{2.6}
\end{equation*}
$$

From (2.5) and (2.6), one gets

$$
\begin{equation*}
\operatorname{ch}\left(E, \mathbf{A}_{0}\right)-\operatorname{ch}\left(E, \mathbf{A}_{T}\right)=d \gamma(T) \tag{2.7}
\end{equation*}
$$

Set $M_{1}=M \backslash \bigcup_{p \in Z(v)} B_{p}(1)$.
Since $V$ is invertible on $M_{1}$, by proceeding as in [5, §4], one sees that the following identity holds uniformly on $M_{1}$,

$$
\begin{equation*}
\lim _{T \rightarrow+\infty} \operatorname{ch}\left(E, \mathbf{A}_{T}\right)=0 \tag{2.8}
\end{equation*}
$$

Lemma 2.1. The following identity holds,

$$
\begin{equation*}
\left\langle\operatorname{ch}\left(E_{+}\right)-\operatorname{ch}\left(E_{-}\right),[M]\right\rangle=-\sum_{p \in Z(v)} \lim _{T \rightarrow+\infty} \int_{\partial B_{p}(1)} \gamma(T) . \tag{2.9}
\end{equation*}
$$

Proof. Since by our choice the connections $\nabla^{E_{ \pm}}$are the trivial connections when restricted to $\bigcup_{p \in Z(v)} B_{p}(1)$, one has

$$
\begin{equation*}
\left\langle\operatorname{ch}\left(E_{+}\right)-\operatorname{ch}\left(E_{-}\right),[M]\right\rangle=\int_{M} \operatorname{ch}\left(E, \mathbf{A}_{0}\right)=\varphi \int_{M} \operatorname{tr}_{s}\left[e^{-\left(\nabla^{E}\right)^{2}}\right]=\varphi \int_{M_{1}} \operatorname{tr}_{s}\left[e^{-\left(\nabla^{E}\right)^{2}}\right] . \tag{2.10}
\end{equation*}
$$

By (2.7), (2.8) and (2.10), we have

$$
\begin{aligned}
\left\langle\operatorname{ch}\left(E_{+}\right)-\operatorname{ch}\left(E_{-}\right),[M]\right\rangle & =\lim _{T \rightarrow+\infty}\left(\int_{M_{1}} \operatorname{ch}\left(E, \mathbf{A}_{0}\right)-\int_{M_{1}} \operatorname{ch}\left(E, \mathbf{A}_{T}\right)\right) \\
& =\lim _{T \rightarrow+\infty} \int_{M_{1}} d \gamma(T)=\lim _{T \rightarrow+\infty} \int_{\partial M_{1}} \gamma(T) \\
& =-\sum_{p \in Z(v)} \lim _{T \rightarrow+\infty} \int_{\partial B_{p}(1)} \gamma(T),
\end{aligned}
$$

where the last equality comes from the orientation consideration. Q.E.D.
Recall that the map $v_{p}$ is the restriction of $v$ on $\partial B_{p}(1)$ (cf. (1.1)).
Lemma 2.2. For any $p \in Z_{v}$, the following identity holds,

$$
\begin{equation*}
\lim _{T \rightarrow+\infty} \int_{\partial B_{p}(1)} \gamma(T)=(-1)^{n} \operatorname{deg}\left(v_{p}\right) \tag{2.11}
\end{equation*}
$$

Proof. For any $p \in Z(v)$, since when restricted on the sphere $\partial B_{p}(1)$, the homomorphism $v$ has been deformed to be unitary, we get that $v^{*}=v^{-1}$ and $V^{2}$ is the identity map acting on $\left.E\right|_{\partial B_{p}(1)}$. Also, since $\nabla^{E}$ is the trivial connection over $B_{p}(1)$, we will use the simplified notation $d$ for it. By (2.3), one has on $B_{p}(1)$ that

$$
A_{t}=d+t V, \quad A_{t}^{2}=d^{2}+t[d, V]+t^{2} V^{2}=t^{2} \operatorname{Id}_{E}+t d V
$$

One then deduces that

$$
\begin{aligned}
& \int_{\partial B_{p}(1)} \gamma(T)=\frac{1}{\sqrt{2 \pi \sqrt{-1}}} \varphi \int_{\partial B_{p}(1)} \int_{0}^{T} \operatorname{tr}_{s}\left[V e^{-\mathbf{A}_{t}^{2}}\right] d t \\
= & \frac{1}{\sqrt{2 \pi \sqrt{-1}}} \varphi \int_{\partial B_{p}(1)} \int_{0}^{T} e^{-t^{2}} \operatorname{tr}_{s}\left[V e^{-t d V}\right] d t \\
= & \frac{1}{(2 \pi \sqrt{-1})^{n}} \frac{-1}{(2 n-1)!} \int_{0}^{T} t^{2 n-1} e^{-t^{2}} d t \int_{\partial B_{p}(1)}\left(\operatorname{tr}_{E_{+}}\left[v^{*} d v\left(d v^{*} d v\right)^{n-1}\right]-\operatorname{tr}_{E_{-}}\left[v d v^{*}\left(d v d v^{*}\right)^{n-1}\right]\right) \\
= & \frac{1}{(2 \pi \sqrt{-1})^{n}} \frac{2(-1)^{n}}{(2 n-1)!} \int_{0}^{T} t^{2 n-1} e^{-t^{2}} d t \int_{\partial B_{p}(1)} \operatorname{tr}_{E_{+}}\left[\left(v^{-1} d v\right)^{2 n-1}\right] .
\end{aligned}
$$

Hence,

$$
\begin{aligned}
\lim _{T \rightarrow+\infty} \int_{\partial B_{\epsilon}(p)} \gamma(T) & =\frac{1}{(2 \pi \sqrt{-1})^{n}} \frac{2(-1)^{n}}{(2 n-1)!} \int_{0}^{+\infty} t^{2 n-1} e^{-t^{2}} d t \int_{\partial B_{p}(1)} \operatorname{tr}_{E_{+}}\left[\left(v^{-1} d v\right)^{2 n-1}\right] \\
& =\frac{1}{(2 \pi \sqrt{-1})^{n}} \frac{(-1)^{n}(n-1)!}{(2 n-1)!} \int_{\partial B_{p}(1)} \operatorname{tr}_{E_{+}}\left[\left(v^{-1} d v\right)^{2 n-1}\right] \\
& =(-1)^{n} \operatorname{deg}\left(v_{p}\right),
\end{aligned}
$$

where one compares with [2, Propositions 1.2 and 1.4] for the last equality. Q.E.D.

From Lemmas 2.1 and 2.2, one gets Theorem 1.1. Q.E.D.
We conclude this section with the following result which is complementary to Theorem 1.1.
Lemma 2.3. Under the Basic Assumption 1.0, for any closed form $\alpha \in \Omega^{*}(M)$ without degree zero component, one has

$$
\begin{equation*}
\left\langle[\alpha]\left(\operatorname{ch}\left(E_{+}\right)-\operatorname{ch}\left(E_{-}\right)\right),[M]\right\rangle=0 \tag{2.12}
\end{equation*}
$$

where $[\alpha] \in H^{*}(M, \mathbf{C})$ is the de Rham cohomology class induced by $\alpha$.
Proof. By the Poincaré lemma (cf. [1]), as $\alpha$ is closed and contains no zero degree component, on each $B_{p}(1), p \in Z(v)$, there exists a form $\beta_{p}$ such that $\alpha=d \beta_{p}$ on an open neighborhood of $B_{p}(1)$.

By partition of unity, one then constructs a differential form $\beta$ on $M$ such that $\beta=\beta_{p}$ on each $B_{p}(1), p \in Z(v)$. Then,

$$
\begin{equation*}
\alpha-d \beta=0 \tag{2.13}
\end{equation*}
$$

on $\cup_{p \in Z(v)} B_{p}(1)=M \backslash M_{1}$.
On the other hand, by (2.4) and (2.5) one knows that for any $t \geq 0$, one has

$$
\begin{equation*}
\left\langle[\alpha]\left(\operatorname{ch}\left(E_{+}\right)-\operatorname{ch}\left(E_{-}\right)\right),[M]\right\rangle=\int_{M}(\alpha-d \beta) \varphi \operatorname{tr}_{s}\left[e^{-\mathbf{A}_{t}^{2}}\right] . \tag{2.14}
\end{equation*}
$$

From (2.8), (2.13) and (2.14), and by taking $t \rightarrow+\infty$, one gets (2.12). Q.E.D.

## 3 A Poincaré-Hopf formula for complex vector fields

Let $M$ be a closed and oriented manifold of dimension $2 n$. Let $g^{T M}$ be a Riemannian metric on $T M$. Let $T_{\mathbf{C}} M=T M \otimes \mathbf{C}$ be the complexification of $T M$. Then $g^{T M}$ extends to a symmetric bilinear form $h^{T_{\mathbf{C}} M}$ on $T_{\mathbf{C}} M$.

Let $K=\xi+\sqrt{-1} \eta \in \Gamma\left(T_{\mathbf{C}} M\right)$ be a complex vector field on $M$, with $\xi, \eta \in \Gamma(T M)$. Then one has

$$
\begin{equation*}
h^{T_{\mathbf{C}} M}(K, K)=|\xi|_{g^{T M}}^{2}-|\eta|_{g^{T M}}^{2}+\sqrt{-1}\langle\xi, \eta\rangle_{g^{T M}} . \tag{3.1}
\end{equation*}
$$

Let $Z_{K}$ be the zero set of $h^{T_{\mathbf{C}} M}(K, K)$, that is,

$$
\begin{equation*}
Z_{K}=\left\{x \in M: h^{T_{\mathbf{C}} M}(K(x), K(x))=0\right\} . \tag{3.2}
\end{equation*}
$$

In the rest of this section, we make the following assumption.
Basic Assumption 3.0. The set $Z_{K}$ consists of a finite number of points.
Let $a_{0}>0$ be the injectivity radius of $g^{T M}$. Let $0<\epsilon<\frac{a_{0}}{2}$.
For any $p \in Z_{K}$, let $B_{p}(\epsilon)=\left\{x \in M: d^{g^{T M}}(x, p) \leq \epsilon\right\}$ be the Riemannian ball centered at $p$. We may take $\epsilon$ small enough so that each $B_{p}(\epsilon)$ does not contain points in $Z_{K} \backslash\{p\}$.

Let $S\left(T B_{p}(\epsilon)\right)=S_{+}\left(T B_{p}(\epsilon)\right) \oplus S_{-}\left(T B_{p}(\epsilon)\right)$ be the Hermitian bundle of spinors associated with $\left(T B_{p}(\epsilon),\left.g^{T M}\right|_{B_{p}(\epsilon)}\right)$. Let $\tau$ be the the involution on $S\left(T B_{p}(\epsilon)\right)$ such that $\left.\tau\right|_{S_{ \pm}\left(T B_{p}(\epsilon)\right)}=$ $\pm\left.\mathrm{Id}\right|_{S_{ \pm}\left(T B_{p}(\epsilon)\right)}$. Let $c(\cdot)$ denote the Clifford action on $S\left(T B_{p}(\epsilon)\right) .{ }^{2}$

[^1]Let $v_{K}(p): \Gamma\left(S_{+}\left(T B_{p}(\epsilon)\right)\right) \rightarrow \Gamma\left(S_{-}\left(T B_{p}(\epsilon)\right)\right)$ be defined by

$$
\begin{equation*}
v_{K}(p)=\tau c(\xi)+\sqrt{-1} c(\eta) \tag{3.3}
\end{equation*}
$$

Then one can prove (see Lemma 3.2 below) that the restriction of $v_{K}(p)$ on the sphere $\partial B_{p}(\epsilon)$ is invertible. Thus it defines an integer $\operatorname{deg}\left(\left.v_{K}(p)\right|_{\partial B_{p}(\epsilon)}\right) \in \mathbf{Z}=K^{1}\left(\partial B_{p}(\epsilon)\right)$.

We can now state the main result of this section as follows.

Theorem 3.1. Under the Basic Assumption 3.0, (i) If $n \geq 2$, then the following identity holds,

$$
\begin{equation*}
\chi(M)=-\sum_{p \in Z_{K}} \operatorname{deg}\left(\left.v_{K}(p)\right|_{\partial B_{p}(\epsilon)}\right) \tag{3.4}
\end{equation*}
$$

(ii) If $n=1$, set $Z_{K,+}=\left\{x \in Z_{K}: \xi(x), \eta(x)\right.$ form an oriented frame at $\left.x\right\}$, then

$$
\begin{equation*}
\chi(M)=-\sum_{p \in Z_{K} \backslash Z_{K,+}} \operatorname{deg}\left(\left.v_{K}(p)\right|_{\partial B_{p}(\epsilon)}\right) \tag{3.4}
\end{equation*}
$$

Proof. For simplicity, we first assume that $M$ is spin and denote by $S(T M)=S_{+}(T M) \oplus$ $S_{-}(T M)$ the bundle of spinors associated with $\left(T M, g^{T M}\right)$.

Let $v_{K}=\tau c(\xi)+\sqrt{-1} c(\eta): S_{+}(T M) \rightarrow S_{-}(T M)$ be defined similarly as in (3.3), only that now it is defined on the whole manifold $M$.

Let $Z\left(v_{K}\right)$ denote the set of points at which $v_{K}$ is not invertible.
Lemma 3.2. One has, (i) If $n \geq 2$, then $Z\left(v_{K}\right)=Z_{K}$; (ii) If $n=1$, then $Z\left(v_{K}\right)=Z_{K} \backslash Z_{K,+}$. Proof. From (3.1) and (3.2), it is clear that $p \in Z_{K}$ if and only if $|\xi|=|\eta|$ and $\langle\xi, \eta\rangle=0$.

Let $v_{K}^{*}: S_{-}(T M) \rightarrow S_{+}(T M)$ be the adjoint of $v_{K}$ with respect to the natural Hermitian metrics on $S_{ \pm}(T M)$. Set $V_{K}=v_{K}+v_{K}^{*}: S(T M) \rightarrow S(T M)$. Then $v_{K}$ is not invertible if and only if $V_{K}^{2}$ is not strictly positive.

Clearly,

$$
\begin{equation*}
V_{K}=\tau c(\xi)+\sqrt{-1} c(\eta): S(T M) \rightarrow S(T M) \tag{3.5}
\end{equation*}
$$

From (3.5), one finds

$$
\begin{equation*}
V_{K}^{2}=|\xi|^{2}+|\eta|^{2}+\sqrt{-1} \tau(c(\xi) c(\eta)-c(\eta) c(\xi)) \tag{3.6}
\end{equation*}
$$

Now if at some $x \in M,|\xi|=|\eta|$ and $\langle\xi, \eta\rangle=0$, then $V_{K}^{2}=2|\xi|^{2}+2 \sqrt{-1} \tau c(\xi) c(\eta)$ which is clearly seen not invertible if $n \geq 2$ or if $n=1$ but $\xi$ and $\eta$ do not form an oriented frame at $x$. ${ }^{3}$

Thus, one gets $Z_{K} \backslash Z_{K,+} \subset Z\left(v_{K}\right)$.
On the other hand, observe that if $|\xi| \neq|\eta|$, then $|\xi|^{2}+|\eta|^{2}>2|\xi| \cdot|\eta|$, while it is clear that $2|\xi| \cdot|\eta|+\sqrt{-1} \tau(c(\xi) c(\eta)-c(\eta) c(\xi)) \geq 0$.

Thus if $|\xi(x)| \neq|\eta(x)|$, then $x$ is not in $Z\left(v_{K}\right)$.

[^2]Now if at some $x \in M,|\xi|=|\eta|$ and $\langle\xi, \eta\rangle \neq 0$, one has

$$
\begin{equation*}
c(\xi) c(\eta)-c(\eta) c(\xi)=c(\xi) c\left(\eta-\frac{\langle\eta, \xi\rangle}{|\xi|^{2}} \xi\right)-c\left(\eta-\frac{\langle\eta, \xi\rangle}{|\xi|^{2}} \xi\right) c(\xi), \tag{3.7}
\end{equation*}
$$

with

$$
\begin{equation*}
\left|\eta-\frac{\langle\eta, \xi\rangle}{|\xi|^{2}} \xi\right|<|\eta| . \tag{3.8}
\end{equation*}
$$

From (3.6)-(3.8), one finds that if at some $x \in M,|\xi|=|\eta|$ and $\langle\xi, \eta\rangle \neq 0$, then $V_{K}^{2}>0$.
Thus, $Z\left(v_{K}\right) \subset Z_{K}$. Moreover, if $n=1$, then one verifies directly that $Z\left(v_{K}\right) \subset Z_{K} \backslash Z_{K,+}$. The proof of Lemma 3.2 is completed. Q.E.D.

Back to the proof of Theorem 3.1. By Lemma 3.2, we know that the Basic Assumption 0.1 holds for $v_{K}: S_{+}(T M) \rightarrow S_{-}(T M)$. Thus one may apply Theorem 1.1 to it to get

$$
\begin{equation*}
\left\langle\operatorname{ch}\left(S_{+}(T M)\right)-\operatorname{ch}\left(S_{-}(T M)\right),[M]\right\rangle=(-1)^{n-1} \sum_{p \in Z\left(v_{K}\right)} \operatorname{deg}\left(\left.v_{K}(p)\right|_{\partial B_{p}(\epsilon)}\right) . \tag{3.9}
\end{equation*}
$$

On the other hand, it is standard that (cf. [4])

$$
\begin{equation*}
\left\langle\operatorname{ch}\left(S_{+}(T M)\right)-\operatorname{ch}\left(S_{-}(T M)\right),[M]\right\rangle=(-1)^{n} \chi(M) \tag{3.10}
\end{equation*}
$$

From (3.9) and (3.10), one gets (3.4).
Thus we have proved Theorem 3.1 in the case where $M$ is spin.
For the general case where $M$ need not be spin, we may consider the Signature complex (cf. [4]) associated with $\left(T M, g^{T M}\right)$ instead. Then the same argument above leads to formulas similar to (3.9) and (3.10), with the right hand sides both be multiplied by a factor $2^{n}$, while in the left hand sides the Spin complex be replaced by the Signature complex. Thus one gets again (3.4). We leave the details to the interested reader.

The proof of Theorem 3.1 is completed. Q.E.D.
Remark 3.3. If $Z_{K}=\emptyset$, then one recovers (and at the same time gives a new proof of) the vanishing result of Jacobowitz [3] which has been stated in Proposition 1.2.

Remark 3.4. Theorem 3.1, in its most general form, should be regarded as a geometric result. As a simple amazing consequence (actually a consequence of Proposition 1.2), if $\chi(M) \neq 0$ and $K=\xi+\sqrt{-1} \eta \in \Gamma\left(T_{\mathbf{C}} M\right)$ is nowhere zero over $M$, then for any Riemannian metric $g^{T M}$ on $T M$, there is at least one point $x \in M$, at which one has $|\xi|_{g^{T M}}=|\eta|_{g^{T M}}$ and $\langle\xi, \eta\rangle_{g^{T M}}=0$. Moreover, if $n=1$, then there exists at least two such points. ${ }^{4}$

Remark 3.5. One may also extend Theorem 3.1 to the case where $T M$ is replaced by an arbitrary oriented Euclidean vector bundle. We leave the details to the interested reader.

[^3]Next, we show that Theorem 3.1 is indeed a generalization of the original Poincaré-Hopf index formula (cf. [1, Theorem 11.25]).

To do so, we take $\xi=0$, then $Z_{K}$ is the zero set of $\eta$, which we have assumed to consist of isolated points.

Without loss of generality we also assume that $|\eta|=1$ on each $\partial B_{p}(\epsilon), p \in Z_{K}$.
In view of the last equality in the proof of Lemma 2.2, one has

$$
\begin{equation*}
\operatorname{deg}\left(\left.v_{K}(p)\right|_{\partial B_{p}(\epsilon)}\right)=\frac{1}{(2 \pi \sqrt{-1})^{n}} \frac{(n-1)!}{(2 n-1)!} \int_{\partial B_{p}(\epsilon)} \operatorname{tr}_{S_{+}(T M)}\left[\left(v^{-1} d v\right)^{2 n-1}\right] \tag{3.11}
\end{equation*}
$$

with

$$
\begin{equation*}
v=\sqrt{-1} c\left(\left.\eta\right|_{\partial B_{p}(\epsilon)}\right) \tag{3.12}
\end{equation*}
$$

Let $f_{1}, \cdots, f_{2 n-1}$ be an orthonormal basis of $T\left(\partial B_{p}(\epsilon)\right)$, let $f_{1}^{*}, \cdots, f_{2 n-1}^{*}$ be the metric dual basis of $T^{*}\left(\partial B_{p}(\epsilon)\right)$.

From (3.12), one deduces that (compare with [6, (27)])
$\operatorname{tr}_{S_{+}(T M)}\left[\left(v^{-1} d v\right)^{2 n-1}\right]=-2^{n-1}(2 n-1)!(\sqrt{-1})^{n} f_{1}^{*} \wedge \cdots \wedge f_{2 n-1}^{*} \int^{B} \eta^{*} \wedge\left(\nabla_{f_{1}}^{T M} \eta\right)^{*} \wedge \cdots \wedge\left(\nabla_{f_{2 n-1}}^{T M} \eta\right)^{*}$,
where $\nabla^{T M}$ is the Levi-Civita connection of $g^{T M}$ and where $\int^{B} \eta^{*} \wedge\left(\nabla_{f_{1}}^{T M} \eta\right)^{*} \wedge \cdots \wedge\left(\nabla_{f_{2 n-1}}^{T M} \eta\right)^{*}$ is the function on $\partial B_{p}(\epsilon)$ such that

$$
\begin{equation*}
\eta^{*} \wedge\left(\nabla_{f_{1}}^{T M} \eta\right)^{*} \wedge \cdots \wedge\left(\nabla_{f_{2 n-1}}^{T M} \eta\right)^{*}=\left(\operatorname{dvol}_{g^{T M}}\right) \int^{B} \eta^{*} \wedge\left(\nabla_{f_{1}}^{T M} \eta^{*}\right) \wedge \cdots \wedge\left(\nabla_{f_{2 n-1}}^{T M} \eta\right)^{*} \tag{3.14}
\end{equation*}
$$

on $\left.\Lambda^{2 n}\left(T^{*} M\right)\right|_{\partial B_{p}(\epsilon)}$.
Let $\eta_{p}: \partial B_{p}(\epsilon) \rightarrow S^{2 n-1}(1)$ denote the canonical map induced by $\left.\eta\right|_{\partial B_{p}(\epsilon)}$.
By (3.14), one finds

$$
\begin{equation*}
f_{1}^{*} \wedge \cdots \wedge f_{2 n-1}^{*} \int^{B} \eta^{*} \wedge\left(\nabla_{f_{1}}^{T M} \eta\right)^{*} \wedge \cdots \wedge\left(\nabla_{f_{2 n-1}}^{T M} \eta\right)^{*}=\eta_{p}^{*} \omega, \tag{3.15}
\end{equation*}
$$

where $\omega$ is the volume form on $S^{2 n-1}(1)$.
From (3.11), (3.13) and (3.15), one gets

$$
\begin{align*}
\operatorname{deg}\left(\left.v_{K}(p)\right|_{\partial B_{p}(\epsilon)}\right) & =-\frac{1}{(2 \pi \sqrt{-1})^{n}} \frac{(n-1)!}{(2 n-1)!} 2^{n-1}(2 n-1)!(\sqrt{-1})^{n} \int_{\partial B_{p}(\epsilon)} \eta_{p}^{*} \omega \\
& =\frac{-(n-1)!}{2 \pi^{n}} \int_{\partial B_{p}(\epsilon)} \eta_{p}^{*} \omega=-\operatorname{deg}\left(\eta_{p}\right), \tag{3.16}
\end{align*}
$$

where $\operatorname{deg}\left(\eta_{p}\right)$ denotes the Brouwer degree (cf. [1]) of the map $\eta_{p}: \partial B_{p}(\epsilon) \rightarrow S^{2 n-1}(1)$.
From (3.4) and (3.16), one gets

$$
\chi(M)=\sum_{p \in \text { zero set of } \eta} \operatorname{deg}\left(\eta_{p}\right),
$$

which is exactly the original Poincaré-Hopf index formula (cf. [1, Theorem 11.25]).

Remark 3.6. Continuing Remark 3.4 and assume $n \geq 2$. Let $K=\xi+\sqrt{-1} \eta$ be such that the zero set of $\xi$ is discrete and that $p \in M$ is a zero point of $\xi$ such that $\operatorname{deg}\left(\xi_{p}\right) \neq \chi(M)$, while $\eta$ vanishes on a closed ball of a sufficiently small positive radius around $p$ and is nowhere zero outside this closed ball. ${ }^{5}$ Then according to (3.16), $-\operatorname{deg}\left(v_{K}(p)\right)=\operatorname{deg}\left(\xi_{p}\right) \neq \chi(M)$. Combining this with Theorem 3.1, we see that for any Riemannian metric $g^{T M}$, there is $x \in M$ such that $|\xi|_{g^{T M}}=|\eta|_{g^{T M}} \neq 0$ and $\langle\xi, \eta\rangle_{g^{T M}}=0$. This extends Remark 3.4 to the case where $K=\xi+\sqrt{-1} \eta$ might vanish on $M$.

Now we exhibit an example to illustrate the last line in Remark 3.4.
Example 3.7. Let $S^{2}=\left\{(x, y, z): x^{2}+y^{2}+z^{2}=1\right\}$ be the standard two sphere in the Euclidean space $\mathbf{R}^{3}$. Set $\xi=(-y, x, 0)$ and $\eta=(z, 0,-x)$. Clearly, as $x^{2}+y^{2}+z^{2}=1, \xi+\sqrt{-1} \eta$ is nowhere zero on $S^{2}$. Now $|\xi|=|\eta|$ together with $\langle\xi, \eta\rangle=0$ imply that $x= \pm 1, y=z=0$. Thus, $Z_{K}$ consists of two points $p=(1,0,0), q=(-1,0,0)$. One then verifies that at $q \in S^{2}$, $\xi=(0,-1,0)$ and $\eta=(0,0,1)$ form an oriented frame of $T_{p} S^{2}$. Thus, by (3.4), one sees that the degree at $p$ equals to -2 , as the Euler number of $S^{2}$ is 2 .

Finally, with the help of Example 3.7, we exhibit an application of Theorem 1.1 in the higher dimensional case.

Example 3.8. We take a product $M=S^{2} \times \cdots \times S^{2}$ with $m \geq 2$ copies of $S^{2}$. We use a subscript to denote the corresponding factor of $S^{2}$. So now let $\xi_{i}, \eta_{i}, 1 \leq i \leq m$, be the vector fields constructed in Example 3.7 on the $i$-th factor $S^{2}$ (denoted by $S_{i}^{2}$ ). Let $v_{K, i}$ be the lifting to $M$ of the corresponding map defined as in the proof of Theorem 3.1 on $S_{i}^{2}$. Then each $v_{K, i}$ maps $\Gamma\left(S_{+}(T M)\right)$ to $\Gamma\left(S_{-}(T M)\right)$. Set $v_{K}=\sum_{i=1}^{m} v_{K, i}$, then one verifies directly that $v_{K}$ is singular only at the point $\left(p_{1}, \cdots, p_{m}\right) \in S^{2} \times \cdots \times S^{2}$. By combining Theorem 1.1 with (3.9) and (3.10), one then gets that the degree of $v_{K}$ at $\left(p_{1}, \cdots, p_{m}\right)$ equals to $-2^{m}$, as the Euler number of $S^{2} \times \cdots \times S^{2}$ equals to $2^{m}$. Conversely, one can compute the degree at $\left(p_{1}, \cdots, p_{m}\right)$ first, and then get the Euler number of $S^{2} \times \cdots \times S^{2}$ by using Theorem 1.1.

## References

[1] R. Bott and L. Tu, Differential Forms in Algebraic Topology. Springer-Verlag, 1982.
[2] E. Getzler, The odd Chern character in cyclic homology and spectral flow. Topology 32 (1993), 489-507.
[3] H. Jacobowitz, Non-vanishing complex vector fields and the Euler characteristic. Proc. Amer. Math. Soc., 137 (2009), 3163-3165.
[4] H. B. Lawson and M.-L. Michelsohn, Spin Geometry. Princeton University Press, 1989.

[^4][5] D. Quillen, Superconnections and the Chern character. Topology, 24 (1985), 89-95.
[6] W. Zhang, $\eta$-invariants and the Poincaré-Hopf index formula. Geometry and Topology of Manifolds X. Eds. W. H. Chen et al. (pp. 336-345). World Scientific, 2000.
H. Feng, School of Mathematics and Statistics, Chongqing University of Technology, Chongqing 400050, PR China
Email: fht@nankai.edu.cn
W. Li, Department of Mathematics, Oklahoma State University, Stillwater, OK 74078-1058, USA
Email: wli@math.okstate.edu
W. Zhang, Chern Institute of Mathematics \& LPMC, Nankai University, Tianjin 300071, PR China
Email: weiping@nankai.edu.cn


[^0]:    *Partially supported by NNSFC, MOEC and NSFC.
    ${ }^{\dagger}$ Partially supported by NNSFC and MOEC.
    ${ }^{1}$ One way to define $\operatorname{deg}\left(v_{p}\right)$ is that $v_{p}$ in (1.1) defines a complex vector bundle $E_{v(p)}$ over a sphere $S^{2 n}\left(v_{p}\right)$ with $\partial B(p)$ as an equator. Then one can define $\operatorname{deg}\left(v_{p}\right)=\left\langle\operatorname{ch}\left(E_{v(p)}\right),\left[S^{2 n}\left(v_{p}\right)\right]\right\rangle$.

[^1]:    ${ }^{2}$ For a thorough treatment of spin geometry involved here, see [4].

[^2]:    ${ }^{3}$ As one verifies in this case that either $\xi=\eta=0$, or $c(\xi)-\sqrt{-1} \tau c(\eta) \neq 0$ while $\left(|\xi|^{2}+\sqrt{-1} \tau c(\xi) c(\eta)\right)(c(\xi)-$ $\sqrt{-1} \tau c(\eta))=0$.

[^3]:    ${ }^{4}$ This is because one can switch $\xi$ and $\eta$.

[^4]:    ${ }^{5}$ The existence of such a vector field is clear, as according to a famous theorem of Hopf, there always exists a vector field on $M$ which vanishes only at $p$.

