

Nutation feature extraction of ballistic missile warhead

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To elucidate nutation signatures, the model of nutation for a ballistic missile warhead is established and formulas of micro-Doppler induced by a target with nutation are derived. Two micro-Doppler features, which are micro-motion period and micro-Doppler bandwidth, are extracted based on high-resolution time-frequency transform. The validity of the features extracted in this reported work is verified by computer simulations even with low signal-to-noise ratio.

Introduction: Neutralising the threat of an incoming ballistic missile is a difficult task, which makes ballistic missile defence (BMD) more and more important. The midcourse phase which is the longest phase of flight, has a good chance of recognising and intercepting [1]. However, warheads flying outside the atmosphere in the midcourse phase are usually accompanied by decoys and debris moving with the same velocity, which presents a great challenge for ballistic missile warhead detection and recognition [2]. Since warheads are usually spin-stabilised, nutation, which is a special signature of a spinning conic warhead, will occur if there is a latitudinal disturbance, which is generally unavoidable. Nutation may induce additional frequency modulations on the returned radar signal, which generate sidebands about the warhead's Doppler frequency, called micro-Doppler frequency. Micro-Doppler can be regarded as a unique signature and provides additional information that is complementary to existing methods [3]. In [4], an algorithm for micro-Doppler feature extraction of a ballistic missile warhead is proposed, however it only deals with spinning-precession without nutation. Reference [5] establishes a mathematical model of nutation for a ballistic missile target and analyses its micro-Doppler by applying time-frequency transform, but does not give an algorithm for feature extraction.

In this Letter, we present a method for micro-Doppler feature extraction of a ballistic missile warhead under the case of nutation motion. Based on high-resolution time-frequency distribution, both the micro-motion period and the micro-Doppler bandwidth are extracted without velocity estimation, which avoids a high computational requirement.

Problem formulation: Without loss of generality, the geometry of radar and a conic-shaped target is depicted in Fig. 1. The origin of the co-ordinate system is the centroid of the conic-shaped target. The target has a coning motion along the axis oz with the unit vector $\mathbf{w}_c = [0, 0, 1]^T$, which is also the axis of the coning. The axis oy which is vertical to the axis oz , lies on the plane where the axis oz and radar line of sight (LOS) oo_1 stay. The right hand co-ordinate system is satisfied between the axis ox and the plane yoz . The target spins around its axis of symmetry oo_2 with an angular velocity Ω_s , while does the coning motion along the axis oz with an angular velocity Ω_c and nutation angle θ , and the axis of symmetry also oscillates up and down with an angular velocity Ω_w and amplitude θ_w simultaneously. The angle of LOS is γ along the unit vector $\mathbf{n} = [0, -\sin \gamma, \cos \gamma]^T$ while the initial nutation angle is θ_0 along the unit vector $\mathbf{w}_s = [0, \sin \theta_0, \cos \theta_0]^T$. The angle between LOS and the axis of symmetry oo_2 is $\varphi (0 \leq \varphi \leq \pi)$, which is also called the attitude angle.

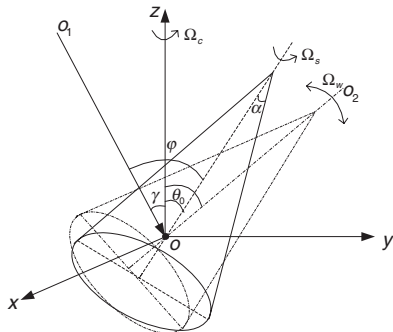


Fig. 1 Geometry of target with nutation motion

According to the Rodrigues's formula, at time t the wobbling rotation matrix is:

$$\mathbf{R}_w = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta_w \sin(\Omega_w t)) & -\sin(\theta_w \sin(\Omega_w t)) \\ 0 & \sin(\theta_w \sin(\Omega_w t)) & \cos(\theta_w \sin(\Omega_w t)) \end{bmatrix} \quad (1)$$

At time t the coning rotation matrix can be written as follows:

$$\mathbf{R}_c = \mathbf{I} + \hat{\mathbf{W}}_c \sin(\Omega_c t) + \hat{\mathbf{W}}_c^2 [1 - \cos(\Omega_c t)] \quad (2)$$

where the skew symmetric matrix is defined by:

$$\hat{\mathbf{W}}_c = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad (3)$$

At time t the unit vector of the axis of symmetry oo_2 becomes

$$\mathbf{r}_1 = \mathbf{R}_c \cdot \mathbf{R}_w \cdot \mathbf{w}_s = [-\sin(\Omega_c t) \sin(\theta_0 - \theta_w \sin(\Omega_w t)), \cos(\Omega_c t) \sin(\theta_0 - \theta_w \sin(\Omega_w t)), \cos(\theta_0 - \theta_w \sin(\Omega_w t))]^T \quad (4)$$

We can write

$$\cos \varphi = \|\mathbf{r}_1 \cdot \mathbf{n}\| / (\|\mathbf{r}_1\| \cdot \|\mathbf{n}\|) = \cos \gamma \cos(\theta_0 - \theta_w \sin(\Omega_w t)) - \sin \gamma \cos(\Omega_c t) \sin(\theta_0 - \theta_w \sin(\Omega_w t)) \quad (5)$$

Therefore, the attitude angle can be expressed by

$$\varphi = a \cos(\cos \gamma \cos(\theta_0 - \theta_w \sin(\Omega_w t)) - \sin \gamma \cos(\Omega_c t) \sin(\theta_0 - \theta_w \sin(\Omega_w t))) \quad (6)$$

The backscattered RCS of a cone can be expressed as [6]

$$\sigma(\varphi) = \begin{cases} \lambda L \tan \alpha \tan^2(\varphi - \alpha) / (8\pi \sin \varphi), & \varphi \in (0, \pi), \varphi \neq \pi/2 - \alpha \\ 8\pi L^3 \sin \alpha / (9\lambda \cos^4 \alpha), & \varphi = \pi/2 - \alpha \end{cases} \quad (7)$$

where λ is the wavelength of the electromagnetic wave. By substituting (6) into the above equation, one can obtain

$$\sigma(t) = \begin{cases} \frac{\lambda L \tan \alpha \tan^2(\cos(\cos \gamma \cos(\theta_0 - \theta_w \sin(\Omega_w t)) - \sin \gamma \cos(\Omega_c t) \sin(\theta_0 - \theta_w \sin(\Omega_w t))) - \alpha)}{8\pi \sin(\cos(\cos \gamma \cos(\theta_0 - \theta_w \sin(\Omega_w t)) - \sin \gamma \cos(\Omega_c t) \sin(\theta_0 - \theta_w \sin(\Omega_w t))))} \cdot \frac{\cos(\cos \gamma \cos(\theta_0 - \theta_w \sin(\Omega_w t)) - \sin \gamma \cos(\Omega_c t) \sin(\theta_0 - \theta_w \sin(\Omega_w t)))}{\sin(\theta_0 - \theta_w \sin(\Omega_w t))} \neq 0, \frac{\pi}{2} - \alpha \\ \frac{8\pi L^3 \sin \alpha}{9\lambda \cos^4 \alpha}, & \cos(\cos \gamma \cos(\theta_0 - \theta_w \sin(\Omega_w t)) - \sin \gamma \cos(\Omega_c t) \sin(\theta_0 - \theta_w \sin(\Omega_w t))) = \frac{\pi}{2} - \alpha \end{cases} \quad (8)$$

Assume the initial position of a scatterer P is $\mathbf{r}_0 = (x_0, y_0, z_0)^T$. Then at time t , the location of the scatterer P is at $\mathbf{r}_t = \mathbf{R}_c \cdot \mathbf{R}_w \cdot \mathbf{r}_0$. If the radar transmits a sinusoidal waveform with a carrier frequency f_0 , the baseband of the signal returned from the particle P is

$$s(t) = \sqrt{\sigma(t)} \exp[j4\pi f_0 r(t)/c] = \sqrt{\sigma(t)} \exp[j\Phi(t)] \quad (9)$$

where $\sqrt{\sigma(t)}$ is the amplitude of the returned radar signal, $r(t) = \|\mathbf{R}_0 + \mathbf{v}t + \mathbf{r}_t\|$, $\mathbf{R}_0$ is the range vector of the origin o from radar at initial time, $\mathbf{v}$ is the velocity vector of bulk motion, c is the speed of the electromagnetic wave. The phase of the baseband signal is $\Phi(t) = 4\pi f_0 r(t)/c$.

By taking the time derivative of the phase and subtracting the Doppler shift caused by translation, the micro-Doppler frequency shift induced by nutation is obtained

$$\begin{aligned} f_{mD}(t) &= \frac{2f_0}{c} \left[\frac{d}{dt} (\mathbf{R}_c \cdot \mathbf{R}_w) \cdot (x_0, y_0, z_0)^T \right]^T \cdot \mathbf{n} \\ &= \frac{2f_0}{c} \left[\frac{d\mathbf{R}_c}{dt} \cdot \mathbf{R}_w \cdot (x_0, y_0, z_0)^T \right. \\ &\quad \left. + \mathbf{R}_c \cdot \frac{d\mathbf{R}_w}{dt} \cdot (x_0, y_0, z_0)^T \right]^T \cdot \mathbf{n} \end{aligned} \quad (10)$$

where

$$\begin{aligned}\frac{d\mathbf{R}_c}{dt} &= \Omega_c[\widehat{\mathbf{W}}_c \cos(\Omega_c t) + \widehat{\mathbf{W}}_c^2 \sin(\Omega_c t)] \\ \frac{d\mathbf{R}_w}{dt} &= \Omega_w \theta_w \cos(\Omega_w t) \\ &\times \begin{bmatrix} 0 & 0 & 0 \\ 0 & -\sin(\theta_w \sin(\Omega_w t)) & -\cos(\theta_w \sin(\Omega_w t)) \\ 0 & \cos(\theta_w \sin(\Omega_w t)) & -\sin(\theta_w \sin(\Omega_w t)) \end{bmatrix}\end{aligned}$$

Feature extraction: From (10), we can conclude that the nutation's micro-Doppler is a time-varying periodic curve. The period and bandwidth of the micro-Doppler can be used as features to recognise a ballistic missile warhead, where the nutation period can be estimated via the autocorrelation method, and the micro-Doppler bandwidth which is twice that of the maximum micro-Doppler frequency can be computed as follows:

1. Compute the smooth pseudo Wigner-Ville distribution of the returned radar signal, written as $\text{SPWVD}(\mathbf{M}, \mathbf{N})$, where M and N are time samples and frequency samples;
2. Calculate the maximum greyscales of the frequencies on the time-frequency plane

$$g_{\max}(f) = \max \{ \text{abs}[(\text{SPWVD}(\mathbf{M}, \mathbf{N}))'] \} \quad (11)$$

where $f = -PRF/2, -PRF/2 + PRF/N, \dots, PRF/2 - PRF/N, PRF$ is pulse repetition frequency;

3. Find the frequencies f_l which satisfy $g_{\max}(f) > \text{mean}\{g_{\max}(f)\}$, where $l = 1, 2, \dots, L$, L is the number of frequencies that satisfy the qualification;
4. Search the position l_0 corresponding to the maximum of $g_{\max}(f)$ and translate it into the centre, which can be realised by the Δl -number cyclic shift of $g_{\max}(f)$, where $\Delta l = \lfloor N/2 - l_0 \rfloor$, $\lfloor \cdot \rfloor$ denotes the round-off operation, $g_{\max}(f)$ can be written as $h_{\max}(f)$ after the cyclic shift;
5. Search from $l = 1$ with sort ascending, until $h_{\max}(f_l) < h_{\max}(f_{l+1}) < h_{\max}(f_{l+2})$, then the maximum negative frequency is $f_{\max-} = f_{l+1}$;
6. Search from $l = L$ with degressive order, until $h_{\max}(f_l) < h_{\max}(f_{l-1}) < h_{\max}(f_{l-2})$, then the maximum positive frequency is $f_{\max+} = f_{l-1}$;
7. The bandwidth of the micro-Doppler can be estimated by the following equation:

$$B_{mD} = f_{\max+} - f_{\max-} \quad (12)$$

A conclusion can be drawn that the proposed algorithm extracts features without bulk velocity estimation, which has the property of requiring less computational burden.

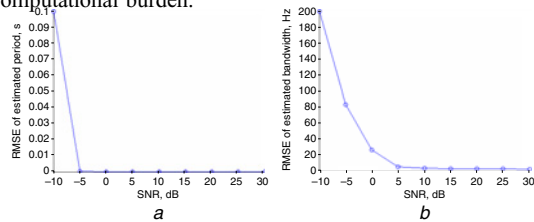


Fig. 2 Variation of micro-Doppler estimation RMSE with SNR

a Micro-Doppler period
b Micro-Doppler bandwidth

Simulation results: Assume that radar operates at 5 GHz and transmits a pulse waveform with a PRF of 1000 Hz. The target whose half cone angle is $\alpha = 10^\circ$, is wobbling along axis oo_2 with an initial nutation angle $\theta_{w0} = 10^\circ$, maximum wobbling angle $\theta_{ww} = 10^\circ$ and angular velocity $\Omega_{ww} = 6/\pi$ rad/s, while it is coning with an angular velocity $\Omega_{wn} = 4\pi$ rad/s. For a spinning invariability conic warhead, scatterers substitute a one by one equivalent, which makes no difference for the returned radar signal. Fig. 2a is the root mean squared error (RMSE) of micro-Doppler period estimation of the autocorrelation method with SNR. The simulation result shows that the estimated period is nearly identical to the theoretic value when $SNR \geq -5$ dB, since the correlation of the signal and randomness of the noise. The RMSE of micro-Doppler bandwidth with SNR is shown in Fig. 2b, which demonstrates that the estimated error of the proposed method is less than 3% when $SNR > 5$ dB. The number of Monte Carlo trials is 100 in all these simulations.

Conclusion: This study establishes the model of nutation for a ballistic missile warhead, and derived mathematical formulas solving micro-Doppler modulations induced by nutation. Based on high-resolution time-frequency transform, the micro-Doppler period and micro-Doppler bandwidth are extracted without velocity estimation, which is computationally efficient. Simulation results confirm that the presented method is effective even with low SNR.

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One or more of the Figures in this Letter are available in colour online.

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