Multi-domain modeling and simulation of proportional solenoid valve

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Abstract: A multi-domain nonlinear dynamic model of a proportional solenoid valve was presented. The electro-magnetic, mechanical and fluid subsystems of the valve were investigated, including their interactions. Governing equations of the valve were derived in the form of nonlinear state equations. By comparing the simulated and measured data, the simulation model is validated with a deviation less than 15%, which can be used for the structural design and control algorithm optimization of proportional solenoid valves.

Key words: fluid mechanics; proportional solenoid valve; dynamic characteristic; multi-domain modeling; simulation

1 Introduction

Proportional solenoid valve is a key control and executing element in industries such as vehicle, construction, aircraft, and mining, whose characteristic is governed by a couple of interacting electro-magnetic, mechanical and flow equations. It has been successfully used for many years due to its high accuracy and robustness which are critical factors for system regular operation. Owing to its importance, a considerable amount of researches on the performance of proportional solenoid valve have been conducted with various methods.

Some dynamic performance simulation research works have been presented to predict the dynamic response of proportional solenoid valve [1-2]. The finite element method (FEM) was usually adopted to analyze the magnetic field and electromagnetic force [3-6]. Hardware-in-loop (HIL) tests of solenoid valve were implemented to view its reliability and performance [7–9]. Efforts on the solenoid driving circuits were made to improve the response performance [10-12]. SETHSON et al [13] proposed a simple model for a 2/2 hydraulic solenoid valve in which the model parameters are tuned to yield good agreement between the measured and the simulated variables. TAGHIZADEH et al [14] developed a nonlinear dynamical model of a pneumatic fast switching valve of which the electro-magnetic, mechanical and fluid subsystems were investigated including their interactions. MAITI et al [15] analyzed the dynamic performance of a proportional pressure relief valve, in which the simulation results were obtained using the MATLAB- SIMULINK environment. DASGUPTA and WATTON [16] proposed a nonlinear dynamic model of a proportional solenoid controlled piloted relief valve through Bondgraph simulation technique, in which higher order dynamics was predicted. WANG et al [17] studied the modeling of the fluid and mechanical subsystems of hydraulic valves using the software SimulationX. However, very few authors have treated the modeling of proportional solenoid valve on a system level to analyze its dynamic characteristics.

As a system, a proportional solenoid valve involves several physical domains and its basic task is to transfer energy among electrical, magnetic, mechanical and hydraulic domains. Various non-linear effects such as flow forces, valve dynamics, orifice flow and chamber fluid compressibility should be taken into account in order to analyze its dynamic characteristics accurately. Therefore, modeling the proportional solenoid valve accurately is a challenging work. In this work, the modeling and identification of a proportional solenoid valve are presented.

2 Structure of valve

Figure 1 shows a simplified schematic view of a proportional solenoid valve (open state) on an 8-speed automatic transmission under study. The solenoid valve is mainly composed of a coil, a core spool and a return spring. Through the moving of the core spool, the valve opens and closes. When the solenoid is energized by DC voltage, the resultant magnetic force displaces the

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Fig.1 Schematic diagram of proportional solenoid valve (open state)

moving spool against the return spring, the valve opens and a flow cross section develops through an orifice.

As shown in Fig.1, a proportional solenoid valve can be decomposed into three subsystems: the electro-magnetic subsystem, mechanical subsystem and fluid subsystem. The electro-magnetic subsystem consists of an electrical and magnetic circuit included in a solenoid. The mechanical subsystem consists of a mass, spring and damper under the effect of magnetic force and hydraulic pressure. These two subsystems behave dynamically and have transient response. The fluid subsystem is the valve orifice, through which the fluid flow is controlled by moving the spool position.

3 Valve modeling

A proportional solenoid value is a complex whose technical system static and dynamic characteristics are governed by a couple of interacting electro-magnetic, mechanical and flow equations. Different subsystems of a proportional solenoid valve are investigated and the governing equations are derived in the form of nonlinear state equations. Therefore, the solenoid current, the spool position and the spool velocity are selected as state variables.

In developing the valve model, the following assumptions are made:

1) A constant source of supply to the inlet port of the valve is considered;

2) The fluid considered for the analysis has Newtonian characteristics;

3) The spring is assumed to be linear;

4) With the opening of the valve ports, the dynamic flow force acting on the valve spool is neglected and only the steady state flow force is considered;

5) The positional striction of the valve spool has not been accounted.

3.1 Electro-magnetic subsystem

This subsystem is used to transfer the input voltage to an electro-magnetic force on the spool of the valve. The magnetic circuit of the valve can be represented by a circuit of discrete resistances [12]. Therefore, the electromagnetic portion in Fig.1 can be translated into a magnetic circuit which consists of iron resistances (R1, R2 and R4), radial air gaps (R3 and R5), a circular variable air gap with stray field (F1), an invariable stray field (G1), a constant air gap with stray field (G2) and an electromagnetic transformer (W1), as shown in Fig.2.



Fig.2 Equivalent magnetic circuit of proportional solenoid valve

According to Kirchihoff's Voltage Law, the following equation must be held for each loop in the magnetic circuit:

$$\sum V_{\mathrm{m},i} = \sum \Phi_i \cdot R_{\mathrm{m},i} = 0 \tag{1}$$

where $V_{\rm m}$ is the magnetic voltage; $R_{\rm m}$ is the magnetic resistance; *i* is the number of the components of the magnetic circuit.

With the definition of the magnetic flux Φ , the following statement is applied to any connection between these components of the magnetic circuit:

$$\sum \Phi_i = 0 \tag{2}$$

The coupling between electric and magnetic quantities through W1 can be deduced from Maxwell's equation:

$$V_{\rm m, W1} = w \cdot I(t) \tag{3}$$

$$V_{\rm EI} = R_{\rm EI}I(t) + w\frac{\mathrm{d}\Phi_{\rm W1}}{\mathrm{d}t} \tag{4}$$

where w is the number of windings; I(t) is the solenoid current; $V_{\rm EI}$ is the electric voltage; $R_{\rm EI}$ is the electric resistance; t is the time.

The constitutive relation of an iron magnetic resistance is determined by its geometry and iron material properties. The characteristics of the iron material can be described by the relationship between the magnetic field strength H and the magnetic flux density B in the form of a bistable sigmoid [12], which is strongly nonlinear. Here, the hysteretic effect is disregarded. Because of the regular geometries of the iron magnetic resistance (R1, R2 and R4), the magnetic voltage and the magnetic flux of an iron resistance can

be approximated as

$$V_{\rm m, R1} = \int_{W} H(r) dr \approx H \cdot l \quad (\text{for R1, R2, R4})$$
(5)

$$\Phi_{\mathrm{R1}} = \int_{r \in A} B(r) \mathrm{d}r \approx B \cdot A \quad (\text{for R1}, \mathrm{R2}, \mathrm{R4}) \tag{6}$$

where H is the magnetic field strength; B is the magnetic flux density; W is the path along the main field in the iron magnetic resistance from the positive to the negative magnetic pin; A is the magnetic cross-section area with normal in B direction; l is the magnetic length in direction of H.

Thus, the iron resistance (R1, R2 and R4) is derived from

$$R_{\rm m, R1} = \frac{V_{\rm m, R1}}{\Phi_{\rm R1}} = \frac{l}{\mu_0 \cdot \mu_{\rm Rel} \cdot A} \quad \text{(for R1, R2, R4)}$$
(7)

The circular air gap F1 is the interface between magnetic subsystem and mechanical subsystem. The stray field is taken into account in addition to the magnetic flux in the air gap. The magnetic resistance of the air gap depends strongly non-linearly on the geometry. The magnetic flux Φ is calculated from its magnetic resistance including its stray flux:

$$\Phi_{\rm F1} = V_{\rm m, F1} \cdot \left(\frac{1}{R_{\rm m, air}} + \frac{1}{R_{\rm m, str}}\right)$$
(8)

$$R_{\rm m,air} = \frac{\mathrm{d}x}{\mu_0.A} \tag{9}$$

$$R_{\rm m, str} = \frac{2.4 \times \sqrt{\frac{4A}{\pi}} + dx}{\mu_0 (1.51 \times \frac{4A}{\pi} + 0.48 \times dx \cdot \sqrt{\frac{4A}{\pi}})}$$
(10)

where $R_{m,air}$ is the air gap resistance; $R_{m,str}$ is the magnetic resistance of its circular stray field; dx is the air gap length; A is the cross-section area.

So, the magnetic resistance of the circular air gap F1 including the stray flux is resulted from

$$R_{\rm m,F1} = \frac{V_{\rm m,F1}}{\Phi_{\rm F1}} = \frac{R_{\rm m,air} \cdot R_{\rm m,str}}{R_{\rm m,air} + R_{\rm m,str}}$$
(11)

Therefore, the magnetic force $F_{m,F1}$ is resulted from

$$F_{m,F1} = \frac{1}{2} \cdot \Phi_{F1}^2 \cdot \frac{dR_{m,F1}}{dx}$$
(12)

As for the constant air gap G2, the magnetic resistance can also be deduced according to Eqs.(8)–(11). The only difference is that the air gap G2 has a constant length while F1 has a variable length.

The winding space can be modeled by an invariable stray field component G1. The magnetic resistance is derived from

$$R_{\rm m,G1} = \frac{4l}{\mu_0 \cdot \pi \cdot (d_{\rm out}^2 - d_{\rm in}^2)}$$
(13)

where d_{out} is the outer diameter of the winding space, d_{in} is its internal diameter, and l is the winding length.

In the radial air gaps, R3 and R5, the space bearing a magnetic flux forms a hollow cylinder. Then, the magnetic resistance is resulted from

$$R_{\rm m,R3} = \frac{\ln\left(1 + a \cdot \left(1 + \sqrt{\frac{2}{a} + 1}\right)\right)}{2 \cdot \mu_0 \cdot \pi \cdot h} \quad (\text{for R3, R5}) \tag{14}$$

where h is the height of the space of the stray field; the variable a depends on its geometry:

$$a = (d_{\text{out}} - d_{\text{in}})^2 / (2d_{\text{in}}^2)$$

3.2 Fluid subsystem

As for the orifice whose turbulent losses are dominant, its hydraulic resistance is described by its geometry. The volume flow Q_{orifice} at the orifices, as shown in Fig.3, is calculated from

$$Q_{\text{orifice}} = \alpha_{\text{D}} \cdot A_{\text{flow}} \cdot \sqrt{\frac{2 \cdot \Delta p}{\rho}}$$
(15)

where A_{flow} is the cross-sectional area; α_D is the flow-rate coefficient; Δp is the pressure difference between point 1 and point 3, $\Delta p = p_1 - p_3$; ρ is the upstream density of the fluid.





According to Hagen-Poiseuille's law, a laminar pressure loss, Δp_{lam} , is added to the turbulent losses:

$$\Delta p_{\text{lam}} = 128 \cdot v \cdot \rho \cdot \frac{x(t)}{\pi \cdot d^4} \cdot Q_{\text{orifice}}$$
(16)

where ν is the upstream kinematic viscosity, ρ is the upstream density, and *d* is the orifice diameter.

The hydrostatic force F_{hy} on the core spool is obtained from

$$F_{\rm hy} = p_{\rm in} \cdot A_{\rm spool} \tag{17}$$

where p_{in} is the pressure of the valve inlet and A_{spool} is the

cross-sectional area of the core spool.

According to the fundamental momentum equation, the steady-state flow force in the axial direction is equal to the momentum difference between the fluid flow into and out of the control volume. Since the flow velocity at the outlet can be neglected compared with the one at the inlet, the axial flow force on the spool can be simplified to

$$F_{\text{flow}} = \rho \cdot Q_{\text{orifice}} \cdot v_{\text{in}} \cdot \cos\theta \tag{18}$$

where v_{in} is the flow velocity at the inlet and θ is the flow angle depending on the opening state of the spool.

3.3 Mechanical subsystem

The transient response of the mechanical subsystem can be represented by Newton's second law as

$$m_{\text{spool}} \cdot \ddot{x}(t) + C_{\text{d}} \cdot \dot{x}(t) + K \cdot (x(t) + \delta) =$$

$$F_{\text{m}}(t) - F_{\text{fr}}(t) - F_{\text{hy}}(t) - F_{\text{fl}}(t) \qquad (19)$$

where m_{spool} is the mass of the core spool; C_d is the damping coefficient; *K* is the spring stiffness coefficient; *x* is the displacement of the core spool; δ is the spring pre-tension; \dot{x} is the velocity of the core spool; \ddot{x} is the acceleration of the core spool; F_m is the magnetic force (see Eq.(12)); F_{hy} is the hydrostatic force (see Eq.(17)); F_{fr} is the slipping frictional force; F_{fl} is the axial steady-state flow force on the spool due to the momentum change of the fluid flow (see Eq.(18)).

4 Model validation

In order to analyze the dynamic characteristics of the proportional solenoid valve and to verify the validation of the above simulation model, a hydraulic test bench was set up, as shown in Fig.4(a). The inlet orifice was set to be 0.9 mm. The pressure and the volume flow at the orifice inlet were measured by two sensors labeled with P and Q, respectively, as shown in Fig.4(b).

The supply pressure was set to be 550 kPa. By controlling PWM width, the input current was increased from 0 to 1 000 mA monotonously and then went back to 0 mA. Commonly, a dither current is adopted to keep the core spool dithering to speed up its dynamic response. Thus, a 50 Hz dither current was added to the control signal. Amplitudes of the dither current were set according to empirical data. Figure 5 shows the dither signal which was added to the PWM control signal for testing. Automatic transmission fluid ATFSSVIII from Sinopec Group was used in the experiment at the standard evaluating temperature of 40 °C.

5 Simulation results and discussion

Combined with the above electro-magnetic, fluid



Fig.4 Test bench for proportional solenoid valves: (a) Test bench picture; (b) Schematic diagram



Fig.5 Dither control signal added to PWM control signal

and mechanical subsystem models, the complete simulation model of a proportional solenoid valve was built based on the multi-domain simulation software platform SimulationX, whose dynamic characteristics were simulated.

Figure 6 displays the calculated dynamic characteristic of the hydraulic portion of the proportional solenoid valve, which shows the strong nonlinearity of the pressure and the volume flow with respect to the motion of the spool.

Figure 7 shows the calculated hydraulic force, spring force, flow force and magnetic force with respect to the current, respectively. It is obvious that the hydraulic force and the magnetic force influence the motion of the core spool significantly.

Figures 8 and 9 show the hydraulic pressure and the volume flow controlled by the proportional solenoid valve with respect to the control current, respectively. It is obvious that the simulation result approximates the test data very well. Both results show that the pressure and the volume flow have strong nonlinear performance with an obvious hysteretic feature on the ascent stage and descent stage. So, the modeling method of the proportional







Fig.7 Simulated forces on core spool







Fig.9 Measured and simulated volume flow

solenoid valve is capable.

The simulation result has less than 15% deviation with the test result. The main reason is that some of the model parameters are hard or impossible to measure, especially when it comes to magnetic material properties. Even for those parameters which can be directly obtained by measurement, there may be measurement errors.

Based on this model, optimized control algorithm for the PWM width could be developed to obtain the required hydraulic pressure.

6 Conclusions

1) A multi-domain nonlinear dynamic model of a proportional solenoid valve system is developed in form of nonlinear state equations and validated by experiment measurements.

2) This model successfully predicts the dynamic characteristics of the valve and can be used as a powerful computational and simulation tool for valve design and algorithm optimization.

3) Unknown parameter and parameter measurement errors are the main reasons for the simulation deviation.

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