

Journal of Alloys and Compounds 432 (2007) 15-17

Journal of ALLOYS AND COMPOUNDS

www.elsevier.com/locate/jallcom

# Magnetic and transport properties of samples of the approximate composition La<sub>0.9</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> with different cation deficiencies

Jin Fang<sup>a</sup>, Yimin Cui<sup>b,\*</sup>

<sup>a</sup> School of Electrical Engineering, Beijing Jiaotong University, Beijing 100044, PR China <sup>b</sup> Physics Department, School of Science, Beijing University of Aeronautics and Astronautics, Beijing 100083, PR China

> Received 9 May 2006; received in revised form 28 May 2006; accepted 31 May 2006 Available online 11 July 2006

## Abstract

Different cation deficiency doped La<sub>0.85</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> and La<sub>0.9</sub>Ba<sub>0.05</sub>MnO<sub>3</sub> and La<sub>0.9</sub>Ba<sub>0.1</sub>Mn<sub>0.95</sub>O<sub>3</sub> and stoichiometric La<sub>0.9</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> polycrystalline bulk samples were synthesized at 1500 °C, the ac magnetic susceptibility, resistivity and magnetic resistance (MR) of the samples were tested, and it is found that stoichiometric phase La<sub>0.9</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> sample have the highest ferromagnetic transition temperature  $T_c$ , and La<sub>0.9</sub>Ba<sub>0.05</sub>MnO<sub>3</sub> sample have the lowest ferromagnetic transition temperature  $T_c$ . La<sub>0.85</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> sample have the largest magnetic resistance (MR), except La<sub>0.9</sub>Ba<sub>0.05</sub>MnO<sub>3</sub> sample, the other samples have the same MR peak temperature.  $(0 \ 2006 Elsevier B.V. All rights reserved.$ 

PACS: 81.40.Rs; 71.30.+h; 75.30.Vn

Keywords: Manganites; ac susceptibility; Resitivity, Magnetic resistance

#### 1. Introduction

Since the observation of colossal magnetoresistance (CMR) effect in doped manganese oxides  $La_{1-x}A_xMnO_3$  (where A is divalent ions, such as Ca, Sr, and Ba), they have been extensively investigated because of their theoretical interests and potential application [1–5]. Although the double exchange interaction between pairs of Mn<sup>3+</sup> and Mn<sup>4+</sup> ions is responsible for the ferromagnetic and metallic properties in these manganese oxides and pairs of Mn<sup>3+</sup> and Mn<sup>4+</sup> can be controlled by changing the doping level [5–7] or oxygen stoichiometry [8], over and above, these properties are also known to strongly depend on the following: average A site ionic size (influencing bond distance d between Mn and O, and bond angle between Mn and oxygen O) [9], the electronegativity [10] and variance [11] of substituted ions, etc. For the cation deficiency doped manganese oxides La<sub>1-x</sub>MnO<sub>3</sub>, significant differences in the magnetic and electric properties have been reported in literatures [12–16]. In this paper, we will report the electrical transport and magnetic properties of La<sub>0.9</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> with different cation deficiency.

## 2. Experimental

The different deficiency doped bulk ceramic samples of  $La_{0.9}Ba_{0.1}MnO_3$ were synthesized using a conventional solid state reaction method. The stoichiometric amounts of  $La_2O_3$ ,  $BaCO_3$  and  $MnO_2$  were mixed, ground and fired at 1300–1400 °C repeatedly. Finally the obtained powders were pelletized into disc samples, and fired at 1500 °C in air for 20 h followed by furnace cooling, respectively. All the samples were characterized by X-ray diffraction (XRD) with Cu K $\alpha$  radiation. The testing samples were taken from the as-prepared disc samples, respectively, the resistivity was measured by the standard four-probe method. AC magnetic susceptibility was measured at 1 kHz by a lock-in technique. The MR was measured using square-probe array method in the magnetic field of 0–1 T. We defined MR ratio as MR ratio (%) =  $(R_0 - R_H)/R_0$ ; where  $R_0$  is the zero-field resistivity and  $R_H$  is the resistivity when the magnetic field was applied.

## 3. Results and discussion

Fig. 1 shows X-ray diffraction patterns of the four different  $La_{0.9}Ba_{0.1}MnO_3$ ,  $La_{0.85}Ba_{0.1}MnO_3$ ,  $La_{0.9}Ba_{0.05}MnO_3$  and  $La_{0.9}Ba_{0.1}Mn_{0.85}O_3$  samples. The XRD patters indicate that the as-prepared samples have the same perovskite structure with almost the same lattice parameter. Within the XRD sensitivity, no obvious difference is detected.

The temperature-dependent ac susceptibility  $\chi$  is shown in Fig. 2. The ferromagnetic transition for all the samples starts at

<sup>\*</sup> Corresponding author. Tel.: +86 10 82339567.

E-mail address: cym@tsinghua.edu.cn (Y. Cui).

<sup>0925-8388/\$ –</sup> see front matter © 2006 Elsevier B.V. All rights reserved. doi:10.1016/j.jallcom.2006.05.125



Fig. 1. X-ray diffraction patterns of the four different samples: (a)  $La_{0.9}Ba_{0.1}MnO_3,$  (b)  $La_{0.85}Ba_{0.1}MnO_3,$  (c)  $La_{0.9}Ba_{0.05}MnO_3$  and (d)  $La_{0.9}Ba_{0.1}Mn_{0.95}O_3.$ 

almost the same temperature (about 230 K), but the transition width sensitively depends on the different ion deficiency. For the stoichiometric  $La_{0.9}Ba_{0.1}MnO_3$  sample, the transition is the sharpest and the susceptibility amplitude is the largest. And for the  $La_{0.9}Ba_{0.05}MnO_3$  sample, the transition peak temperature is the lowest and the susceptibility amplitude is greatly decreased. The  $La_{0.85}Ba_{0.1}MnO_3$  sample and  $La_{0.9}Ba_{0.1}Mn_{0.95}O_3$  sample have almost the same susceptibility, only below 210 K, the former is larger than the latter. At low temperature, susceptibility reduced and spin-glass-like behaviour happened for all samples, which may be explained by spin-cluster model.

Fig. 3 shows the change in resistivity of different samples with temperature. Every sample exhibits semiconductor resistivity and a inflexion nearly at its respective temperature  $T_{\rm C}$ . Within the whole range of temperature we investigated, the resistivity is sensitively depends on the different ion deficiency. For the stoichiometric La<sub>0.9</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> sample, the resistivity is



Fig. 2. The temperature-dependent ac susceptibility  $\chi$  of the four different samples.



Fig. 3. The temperature-dependent resistivity of the four different samples.

the least. The resistivity of  $La_{0.9}Ba_{0.1}Mn_{0.95}O_3$  sample is slight larger than that of  $La_{0.85}Ba_{0.1}MnO_3$  sample. Corresponding to the susceptibility, the resistivity of the  $La_{0.9}Ba_{0.05}MnO_3$  sample is obvious larger than that of the other samples, especially at low temperature, from  $La_{0.9}Ba_{0.1}MnO_3$  sample to  $La_{0.9}Ba_{0.05}MnO_3$ Sample, the resistivity increased by more than two orders.

The magnetoresistance ratio (MR), defined as  $\nabla R/R_0 = (R_0 - R_H)/R_0$ , where  $R_0$  and  $R_H$  are the resistivity of the zero field and an applied field *H*, respectively, is plotted in Fig. 4. Every sample exhibits a MR peak nearly at its respective temperature  $T_C$ . Calculation show the MR ratios for 1 T applied field are more than 20% for the La<sub>0.85</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> sample and 19% for the stoichiometric La<sub>0.9</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> sample, and almost the same 16% for the other two samples.

Those phenomenon has been traditionally understood on the basis of the double-exchange (DE) model and Jahn-Teller effects [17–20]. When La ions in the LaMnO<sub>3</sub> were replaced by a divalent element Ba, proportional amount of  $Mn^{3+}$ , with electronic configuration  $t_{2g}^3 e_g^1$ , is substituted with  $Mn^{4+}$  ( $t_{2g}^3$ ) and form a mixed-valence state of  $Mn^{3+}$ – $Mn^{4+}$ . The hopping magnitude of  $e_g$  electron between spin aligned  $Mn^{3+}$  and  $Mn^{4+}$  ions via



Fig. 4. The temperature-dependent of the magnetoresistance ratio (MR) of the four different samples.

 $O^{2-}$  rests on the angle  $(\theta_{ij})$  between the neighboring spins, and relates to the applied magnetic field. In a hole doped CMR system, the mixed valence state of  $Mn^{3+}-Mn^{4+}$  is a key component for understanding CMR effect and the transition from a ferromagnetic metal to a paramagnetic semiconductor. At low doped concentration ( $x \le 0.2$ , x is the doped concentration of divalent cation), Alonso et al. [21] propose that  $Mn^{4+}$  holes are localized around bivalent cations forming ferromagnetic clusters, and the undistorted clusters around the bivalent ions are also locally conducting although the spatial localization imposed by the bivalent ions dopants results in clusters that are electrically isolated from each other.

The concentration of  $Mn^{4+}$  holes in La<sub>0.9</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> is 10%. For La<sub>0.85</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> and La<sub>0.9</sub>Ba<sub>0.1</sub>Mn<sub>0.95</sub>O<sub>3</sub>, the cation deficiency ratios are only 5%, respectively, and the concentrations of Mn<sup>4+</sup> holes are 25%. But for La<sub>0.9</sub>Ba<sub>0.05</sub>MnO<sub>3</sub>, the Ba cation deficiency ratios are almost 50%, and the concentration of Mn<sup>4+</sup> holes is about 20%, less than that of the two deficiency samples. The deficiencies may exist on defect or small separation of cation oxides which may cause some disorder. Although the Mn<sup>4+</sup> holes of the three deficiency samples are higher than that of the La<sub>0.9</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> sample, the disorder due to vacancies decreases the Curie temperature. At the same time, compare with the stoichiometric La<sub>0.9</sub>Ba<sub>0.1</sub>MnO<sub>3</sub>, the change of bond angle between Mn and oxygen O in La<sub>0.9</sub>Ba<sub>0.05</sub>MnO<sub>3</sub> is larger than those in La<sub>0.85</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> and La<sub>0.9</sub>Ba<sub>0.1</sub>Mn<sub>0.95</sub>O<sub>3</sub>, so the double change of Mn<sup>3+</sup> and Mn<sup>4+</sup> in La<sub>0.9</sub>Ba<sub>0.05</sub>MnO<sub>3</sub> sample is greatly reduced, and the change of susceptibility, resistivity and MR for La<sub>0.9</sub>Ba<sub>0.05</sub>MnO<sub>3</sub> sample is more remarkable than those of the other two cation deficient samples. As for La<sub>0.85</sub>Ba<sub>0.1</sub>MnO<sub>3</sub> sample, have more Mn cation than the other samples, which may improve the MR, but due to the disorder, the resistivity is still larger than that of the stoichiometric sample.

## 4. Conclusion

In summary, we have investigated the ac magnetic and electric and MR properties of different deficiency doped  $La_{0.9}Ba_{0.1}MnO_3$ , the  $T_C$  and amplitude of ferromagnetic transitions for all cation deficiency samples are lower than that of the stoichiometric sample, at the same time the resistivity for all cation deficiency samples are larger than that of the stoichiometric sample. Especially for the Ba cation deficiency sample, the change of double exchange is largest, the changes of sus-

ceptibility and resistivity are largest correspondingly. In all, the disorder due to the vacancies decreases the Curie temperature although the number of Mn<sup>4+</sup> ions increases in the all cation deficiency samples.

## Acknowledgment

We acknowledge the financial support from National 973 Project of China, National Natural Science Foundation of China.

## References

- R. von Helmot, J. Wecker, B. Hopzapfel, L. Schulz, K. Samwer, Phys. Rev. Lett. 71 (1993) 2331.
- [2] S. Jin, T.H. Tiefel, M. McCormack, R.A. Fastnacht, R. Ramesh, L.H. Chen, Science 264 (1994) 413.
- [3] G.C. Xiong, H.L. Qi Li, S.N. Ju, L. Mao, X.X. Senapati, R.L. Xi, T. Greene, Venkatesan, Appl. Phys. Lett. 66 (1995) 1427.
- [4] A. Urushibara, Y. Moritomo, T. Arima, A. Asamitsu, G. Kido, Y. Tokura, Phys. Rev. B 51 (1995) 14103.
- [5] N. Khare, U.P. Moharil, A.K. Gupta, A.K. Raychaudhuri, S.P. Pai, R. Pinto, Appl. Phys. Lett. 81 (2002) 325.
- [6] M. Ziese, G. Heydon, R. Hohne, P. Esquinazi, J. Dienelt, Appl. Phys. Lett. 74 (1999) 1481.
- [7] K. Steenbeck, T. Eick, K. Kirsch, K. O'Donnell, E. Steinbeiss, Appl. Phys. Lett. 71 (1997) 968.
- [8] H.L. Ju, J. Gopalkrishnan, J.L. Peng, Q. Li, G.C. Xiong, T. Venkatesan, R.L. Greene, Phys. Rev. B 51 (1995) 6143.
- [9] H.Y. Hwang, S.W. Cheong, P.G. Radaelli, M. Marezio, B. Batlogg, Phys. Rev. Lett. 75 (1995) 914.
- [10] J. Fontcuberta, J.L. Garcia-Munoz, M. Suaaidi, B. Martinez, S. Pinol, X. Obradors, J. Appl. Phys. 81 (1997) 5481.
- [11] M. Lide, J. Rodriguez-Martinez, P. Atteld, Phys. Rev. B 54 (1996) R15622.
- [12] G. Dezanneau, O. Isnard, H. Roussel, A. Sin, M. Audier, H. Vincent, Crystal Eng. 5 (2002) 347.
- [13] R. Horyn, E. Bukowska, A. Sikora, J. Alloys Compd. 346 (2002) 107.
- [14] A.A. Bosak, O.Y. Gorbenko, A.R. Kaul, I.E. Graboy, C. Dubourdieu, J.P. Senateur, H.W. Zandbergen, J. Magn. Magn. Matter 211 (2000) 61.
- [15] M. Wolcyrz, R. Horyn, F. Bouree, E. Bukowska, J. Alloys Compd. 353 (2003) 170.
- [16] M. Yahia, H. Batis, Euro. J. Inorg. Chem. (13) (2003) 2486–2494.
- [17] C. Zener, Phys. Rev. 82 (1951) 403.
- [18] P.W. Anderson, H. Hasegawa, Phys. Rev. 100 (1955) 675.
- [19] J.B.A.A. Elemans, B. Van Laar, K.R. Van Der Veen, B.O. Loopstra, J. Solid State Chem. 3 (1971) 238.
- [20] A.J. Millis, B.I. Shaiman, R. Mueller, Phys. Rev. Lett. 77 (1996) 175.
- [21] J. Alonso, E. Herrero, J.M. Gonzailez-Calbet, M. Vallet-Regi, J.L. Martinez, J.M. Rojo, A. Hernando1, Phys. Rev. B 62 (2000) 11328.